

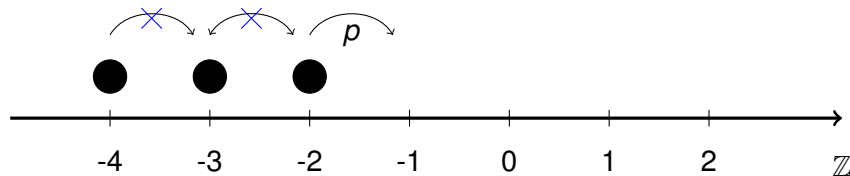
# Cutoff Profile of ASEP on the segment

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joint work with Alexey Bufetov

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Rouen 25.11.2022  
Rencontre de Probabilités

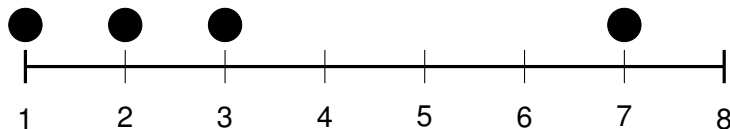
# Asymmetric simple exclusion process (ASEP)



## Dynamics:

- ▶ There is at most one particle per site, at rate one particles jump **one step to the right or left if possible**
- ▶ jumps to the right (resp. left) have probability  $p > 1/2$ , (resp.  $q = 1 - p$ ), so there is a **drift to the right**
- ▶ for  $p = 1$  we obtain the totally ASEP (**TASEP**)
- ▶ Markov process  $(\eta_t, t \geq 0)$  with state space  $\Omega = \{0, 1\}^{\mathbb{Z}}$

# ASEP on the segment



- ▶ put  $k$  particles on the segment  $[1; N] = \{1, \dots, N\}$
- ▶ jumps out of the segment are forbidden
- ▶ Markov chain with stationary measure  $\pi_{N,k}$
- ▶  $P_t^\eta$  = law of the ASEP started from  $\eta$  at time  $t$

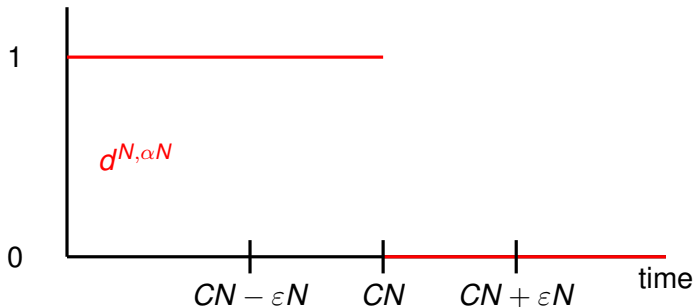
Goal: Understand how ASEP **mixes** to equilibrium w.r.t.

$$d^{N,k}(t) = \max_{\eta} \|P_t^\eta - \pi_{N,k}\|_{\text{TV}} \in [0, 1]$$

as  $N \rightarrow \infty$  and  $k = \alpha N, \alpha \in (0, 1)$ .

# Cutoff

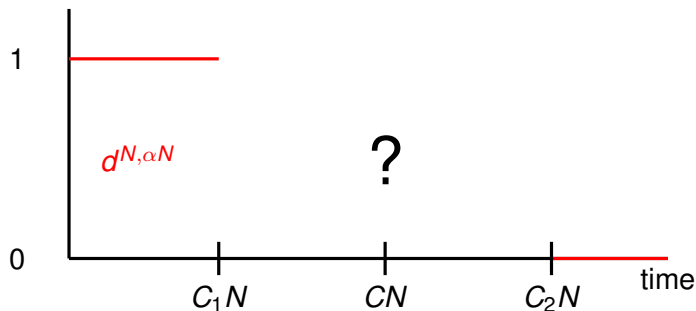
- ▶ Often  $d^{N,\alpha N}$  drops from 1 to 0 abruptly, i.e. **cutoff** holds: At time  $CN \pm \varepsilon N$ ,  $d^{N,\alpha N}$  goes to 1 resp. 0 for any  $\varepsilon > 0$ .



- ▶ simple example of cutoff: ASEP with 1 particle on the segment
- ▶ no cutoff for ASEP with 1 particle on the circle  $\mathbb{Z}_N$

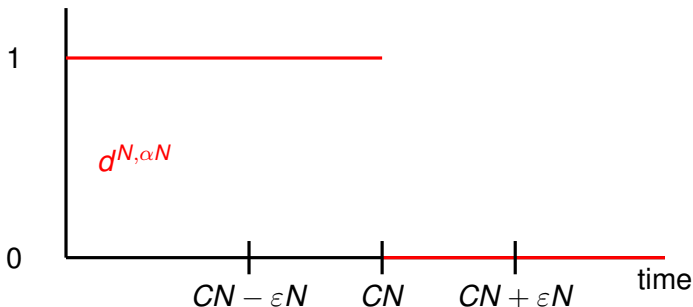
# Precutoff and Cutoff for ASEP

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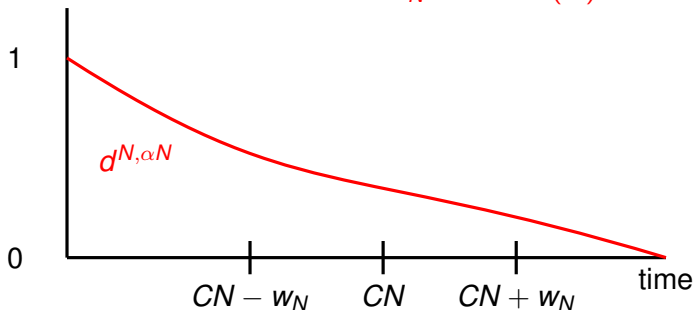
1. [Benjamini et al. '03, Trans. Math. Soc.] showed there is precutoff:
2. [Labbé - Lacoïn '16, Ann. Prob.] showed cutoff at  
$$C = \frac{(\sqrt{\alpha} + \sqrt{1-\alpha})^2}{p-q}$$



Both papers show their result also for the multi-color ASEP

# Cutoff window and profile

When we **zoom in** on the cutoff point, we expect to see a **smooth transition** in a **window  $w_N$**  of size  $o(N)$  :



The function which describes the smooth transition is the **cutoff profile**.

# Cutoff profile

We prove the following:

**Theorem (Bufetov-N. '22, PTRF '22)**

Let  $c \in \mathbb{R}$ , and let  $k_N = \alpha N, \alpha \in (0, 1)$ . Then we have

$$\lim_{N \rightarrow \infty} d^{N, \alpha N} \left( \frac{(\sqrt{\alpha} + \sqrt{1 - \alpha})^2}{p - q} N + \frac{cN^{1/3}}{p - q} \right) = 1 - F_{\text{GUE}}(cf(\alpha)),$$

where  $f(\alpha) = \frac{(\alpha(1 - \alpha))^{1/6}}{(\sqrt{\alpha} + \sqrt{1 - \alpha})^{4/3}}.$

- ▶  $F_{\text{GUE}}$  asymptotic law of the largest eigenvalue of a random matrix from the **G**aussian **U**nitary **E**nsemble
- ▶ Cutoff window is  $N^{1/3}$ , -profile is  $1 - F_{\text{GUE}}$ , Cutoff is reproven
- ▶ replace  $\frac{(\sqrt{\alpha} + \sqrt{1 - \alpha})^2}{p - q} N$  by  $\frac{(\sqrt{k_N} + \sqrt{N - k_N})^2}{p - q}$  when  $\frac{k_N}{N} \rightarrow \alpha$



# Notation and bounds

For brevity, we denote the time point

$$g(k, c) := \frac{(\sqrt{k} + \sqrt{N-k})^2 + cN^{1/3}}{p-q}.$$

and always assume the number of particles  $k_N$  satisfies  $k_N/N \rightarrow \alpha \in (0, 1)$ .

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We will separately show

$$\liminf_{N \rightarrow \infty} d^{N, k_N}(g(k_N, c)) \geq 1 - F_{\text{GUE}}(cf(\alpha)),$$

and

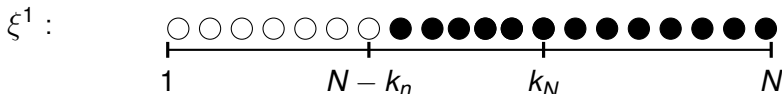
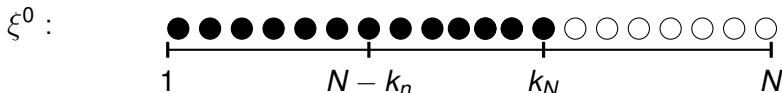
$$\limsup_{N \rightarrow \infty} d^{N, k_N}(g(k_N, c)) \leq 1 - F_{\text{GUE}}(cf(\alpha)), \quad (1)$$

As is often the case, it is harder to show the upper bound (1), so we focus on (1).

# Hitting times I

Consider the configurations for ASEP on the segment

$$\xi^0 = \mathbf{1}_{[1;k_N]}, \quad \xi^1 = \mathbf{1}_{[N-k_N+1;N]}.$$



$\xi^{(0)}$  is intuitively the '**worst**' initial configuration.

Let  $\mathfrak{h}$  the first time that ASEP started from  $\xi^0$  reaches  $\xi^1$  :

$$\mathfrak{h} = \inf\{t : \xi_t^0 = \xi^1\}.$$

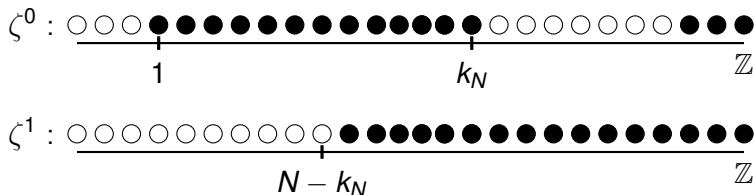
Then we have the inequality

$$d^{N,k_N}(t) \leq \mathbb{P}(\mathfrak{h} > t).$$

## Hitting times II

Consider the configurations for ASEP on  $\mathbb{Z}$

$$\zeta^0 = \mathbf{1}_{[1; k_N]} + \mathbf{1}_{\mathbb{Z}_{>N}}, \quad \zeta^1 = \mathbf{1}_{\mathbb{Z}_{>(N-k_N)}}.$$



Let  $\mathfrak{H}$  the first time that ASEP started from  $\zeta^0$  reaches  $\zeta^1$  :

$$\mathfrak{H} = \inf\{t : \zeta_t^0 = \zeta^1\}.$$

Then we have the inequalities

$$d^{N, k_N}(t) \leq \mathbb{P}(\mathfrak{H} > t) \leq \mathbb{P}(\mathfrak{H} > t)$$

# Upper bound using hitting times

- ▶ The hitting time  $\mathfrak{H}$  was already studied by [Benjamini et al '03].
- ▶ We will **eventually** show

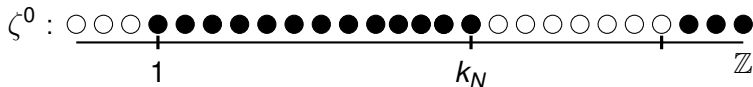
$$\lim_{N \rightarrow \infty} \mathbb{P}(\mathfrak{H} > g(k_N, c)) = 1 - F_{\text{GUE}}(cf(\alpha)).$$

By the inequality  $d^{N, k_N}(t) \leq \mathbb{P}(\mathfrak{H} > t)$ , this will imply

$$\limsup_{N \rightarrow \infty} d^{N, k_N}(g(k_N, c)) \leq 1 - F_{\text{GUE}}(cf(\alpha)).$$

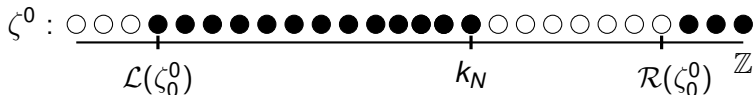
# Leftmost particle / rightmost hole

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## Leftmost particle / rightmost hole

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- ▶ To understand the hitting time  $\mathfrak{H}$ , we track the leftmost particle/rightmost hole of  $\zeta_t^0$ :

$$\mathcal{L}(\zeta_t^0) = \min\{i \in \mathbb{Z} : \zeta_t^0(i) = 1\} \quad \mathcal{R}(\zeta_t^0) = \max\{i \in \mathbb{Z} : \zeta_t^0(i) = 0\}.$$

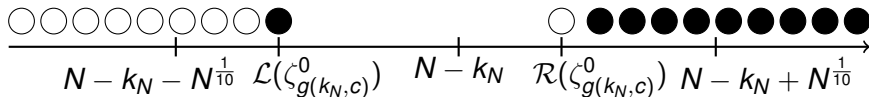
Note that deterministically for all  $t$

$$\mathcal{L}(\zeta_t^0) - 1 \leq N - k_N \leq \mathcal{R}(\zeta_t^0)$$

and  $\mathfrak{H}$  is precisely the first time that

$$\mathcal{L}(\zeta_{\mathfrak{H}}^0) - 1 = N - k_N = \mathcal{R}(\zeta_{\mathfrak{H}}^0)$$

Let  $B_N(c)$  be the event that  $\zeta_{g(k_N, c)}^0$  looks like this:

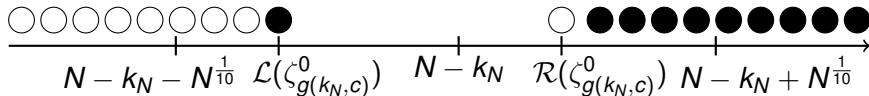


- ▶ when  $B_N(c)$  happens,  $\mathfrak{H}$  cannot be much bigger than  $g(k_N, c)$ , in particular we can show

$$\lim_{N \rightarrow \infty} \mathbb{P}(B_N(c)) = \lim_{N \rightarrow \infty} \mathbb{P}(\mathfrak{H} < g(k_N, c))$$



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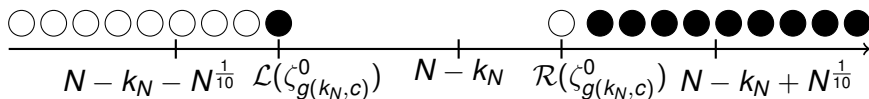
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We want to prove

$$\lim_{N \rightarrow \infty} \mathbb{P}(B_N(c)) = F_{\text{GUE}}(cf(\alpha)).$$

- ▶ This will imply  $\lim_{N \rightarrow \infty} \mathbb{P}(\mathfrak{H} > g(k_N, c)) = 1 - F_{\text{GUE}}(cf(\alpha))$ , and thus yield the upper bound for  $d^{N, k_N}(g(k_N, c))$

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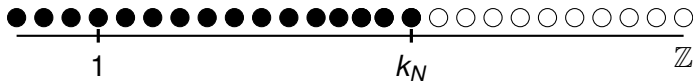
We want to prove

$$\lim_{N \rightarrow \infty} \mathbb{P}(B_N(c)) = F_{\text{GUE}}(cf(\alpha)).$$

Main sources of this convergence are :

- ▶ non-standard CLT for ASEP
- ▶ algebraic identities for multi-color ASEP
- ▶ couplings to compare different ASEPs

# Non-standard CLT for ASEP



We start ASEP on  $\mathbb{Z}$  from  $\mathbf{1}_{\mathbb{Z}_{\leq k_N}}$ , ("step initial data") and denote

$x_{k_N}(t)$  = position at time  $t$  of the particle that started in 1.

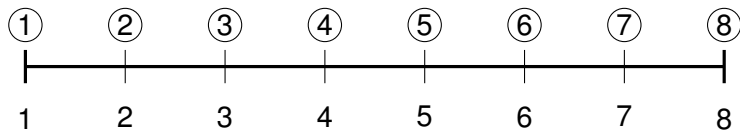
Theorem (Corollary of Johansson '00 (for  $p=1$ ),  
Tracy-Widom '09)

We have for  $k_N$  with  $k_N/N \rightarrow \alpha \in (0, 1)$  that

$$\lim_{N \rightarrow \infty} \mathbb{P}(x_{k_N}(g(k_N, c)) \leq N - k_N) = 1 - F_{\text{GUE}}(cf(\alpha)),$$

where  $f(\alpha) = \frac{(\alpha(1-\alpha))^{1/6}}{(\sqrt{\alpha} + \sqrt{1-\alpha})^{4/3}}.$

# Multi-color ASEP

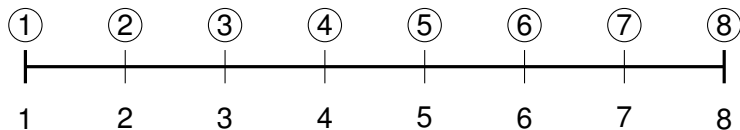


- ▶ we have  $N$  particles on  $[1; N]$  with colors  $1, \dots, N$
- ▶ each particle moves as in ASEP, but jumps to sites occupied by **smaller color** are impossible
- ▶ encoded by a permutation  $\pi$  mapping **positions to colors**
- ▶ invariant measure is the Mallows measure

$$\mathcal{M}(\pi) = \frac{(p/q)^{\#\text{inv}(\pi)}}{Z_N},$$

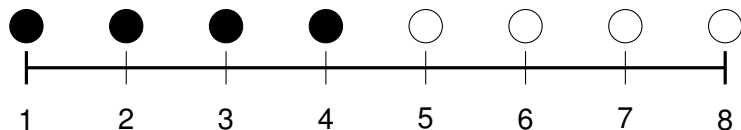
where  $\text{inv}(\pi)$  are the inversions of  $\pi$ , for  $q = 0$ ,  $\mathcal{M}$  is the Dirac measure on  $\pi(i) = N - i + 1$

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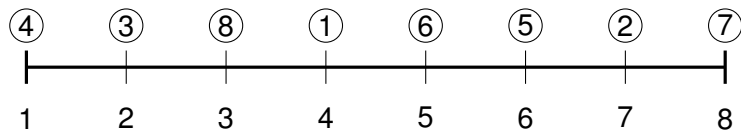
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- ▶ encoded by a permutation  $\pi$  mapping **positions to colors**
- ▶ when we project down and only distinguish between particles of color smaller equal  $k$  and greater than  $k$ , we recover ASEP

# Bringing into equilibrium

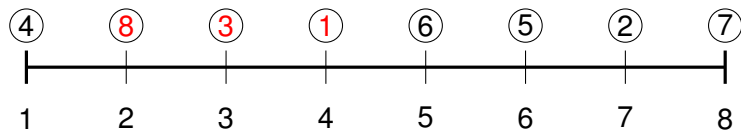
- ▶ Consider a permutation of  $[a; b]$  and let  $[c; d] \subseteq [a; b]$  :
- ▶ Example :  $[a; b] = [1; 8]$  and  $[c; d] = [2; 4]$



- ▶ for TASEP to **bring into equilibrium**  $[c; d]$  means that we order the colors in  $[c; d]$  in decreasing order
- ▶ more generally, it means to distribute the colors in  $[c; d]$  according to the Mallows measure

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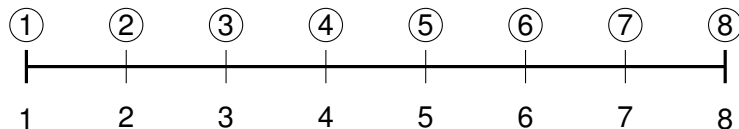


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# Color-position symmetry

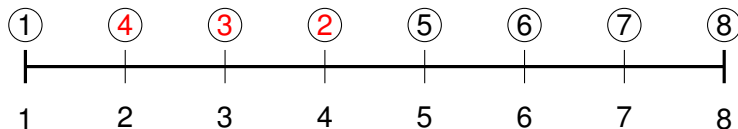
- ▶ we start with the identity permutation



- ▶ bring into equilibrium  $[c_1; d_1]$ , then  $[c_2; d_2]$
- ▶ here:  $[c_1; d_1] = [2; 4]$  and  $[c_2; d_2] = [4; 7]$

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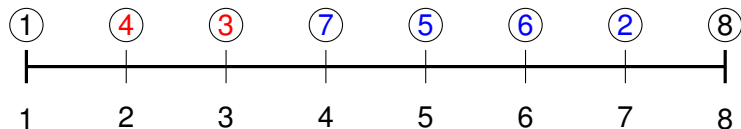
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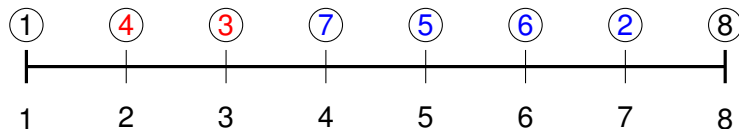
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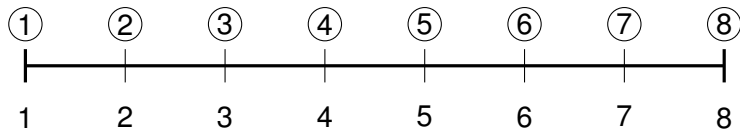
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- ▶ run the process up to time  $t$
- ▶ this results in a permutation  $\pi_t$

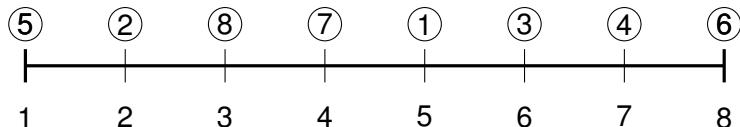
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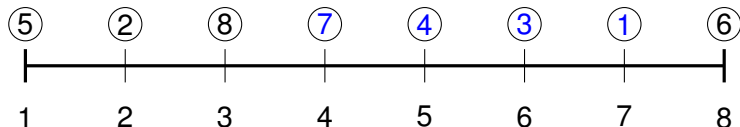
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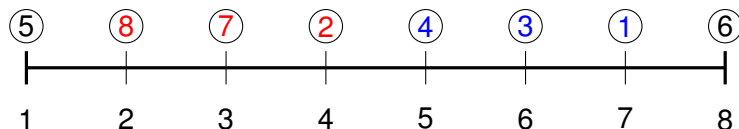
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- ▶ run the process up to time  $t$
- ▶ bring into equilibrium  $[c_2; d_2] = [4; 7]$
- ▶ bring into equilibrium  $[c_1; d_1] = [2; 4]$
- ▶ this results in a permutation  $\hat{\pi}_t$



# Color-position symmetry

We have the following unintuitive identity:

## Proposition (Bufetov-N. '22)

*The permutations  $\pi_t^{-1}$  and  $\hat{\pi}_t$  are equal in law.*

- ▶ source of this identity is purely algebraic (using Hecke algebras, cf. [Bufetov '21+], also [Borodin-Bufetov '20])
- ▶ since  $\pi_t$  maps positions to colors, whereas  $\hat{\pi}_t^{-1}$  maps colors to positions, this is a color-position symmetry

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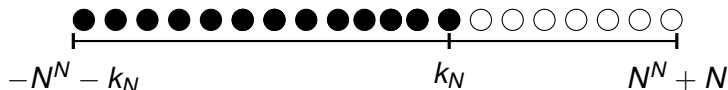
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Main application:

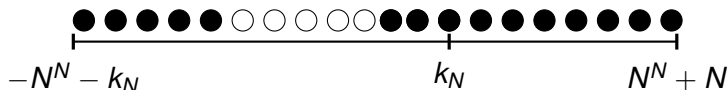
- ▶  $(\pi_t, t \geq 0)$  starts from a 'complicated' permutation, which can be projected down to (a perturbation of)  
 $\zeta^0 = \mathbf{1}_{[1;k_N]} + \mathbf{1}_{\mathbb{Z}_{>N}}$
- ▶  $(\hat{\pi}_t, t \geq 0)$  starts from the identity, which can be projected down to step initial data (CLT available)

# The permutation $\pi_0$

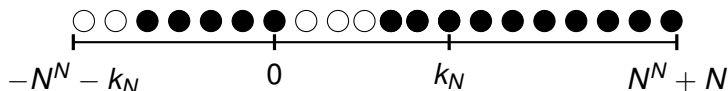
Consider ASEP on a very large segment with this step initial data:



We bring into equilibrium  $[-N^N; N^N + N]$ :



We bring into equilibrium  $[-N^N - k_N; 0]$ :

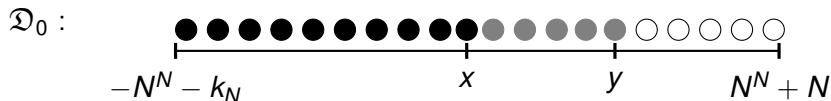


This is very close to a shift of  $\zeta^0$ , we call it  $\hat{\zeta}^0$

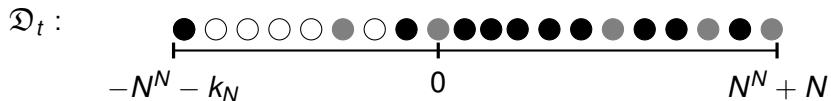
# The permutation $\hat{\pi}_0$

- ▶  $x, y$  free parameters
- ▶ in the identity, we project down on colors  $\leq x$  (black), colors in  $(x, y]$  (grey) and colors  $> y$  (white)

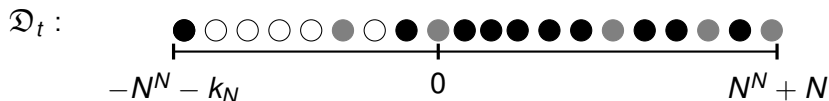
Our initial configuration thus is



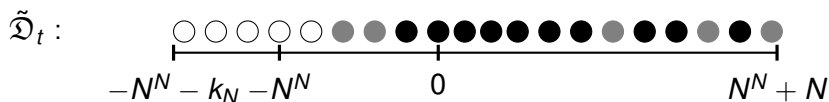
After time  $t$ , the process may look like this:



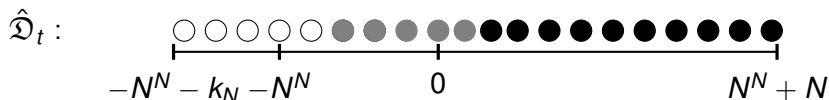
# Bringing into equilibrium



In  $\mathfrak{D}_t$  we bring into equilibrium  $[-N^N - k_N; 0]$  :



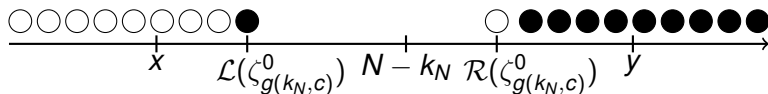
Afterwards, we bring into equilibrium  $[-N^N; N^N + N]$  :



# Reminders

$$x := N - k_N - N^{1/10} \quad y := N - k_N + N^{1/10}$$

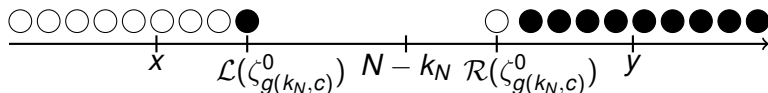
Recall  $B_N(c)$  be the event that  $\zeta_{g(k_N, c)}^0$  looks like this:



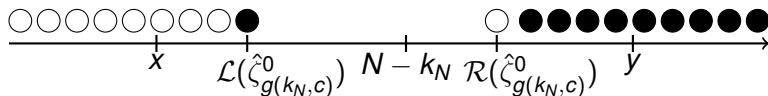
# Reminders

$$x := N - k_N - N^{1/10} \quad y := N - k_N + N^{1/10}$$

Recall  $B_N(c)$  be the event that  $\zeta_{g(k_N, c)}^0$  looks like this:



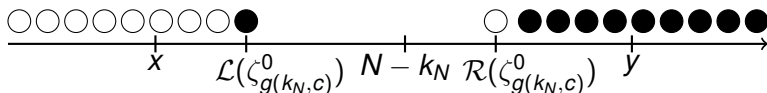
We call  $\hat{B}_N(c)$  be the event that  $\hat{\zeta}_{g(k_N, c)}^0$  looks the same:



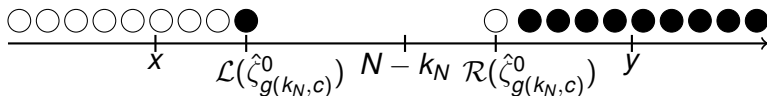
# Reminders

$$x := N - k_N - N^{1/10} \quad y := N - k_N + N^{1/10}$$

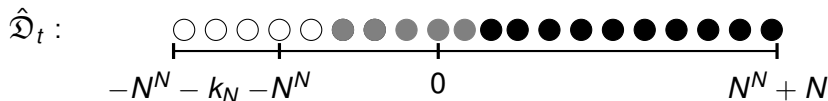
Recall  $B_N(c)$  be the event that  $\zeta_{g(k_N, c)}^0$  looks like this:



We call  $\hat{B}_N(c)$  be the event that  $\hat{\zeta}_{g(k_N, c)}^0$  looks the same:



Recall  $\hat{\mathfrak{D}}_t$  from the previous slide:





As corollary of the color-position symmetry, we have:

### Proposition (Bufetov -N.)

*We have*

$$\mathbb{P}(\hat{B}_N(c)) = \mathbb{P} \left( \begin{array}{l} \text{all black particles in } \hat{\mathfrak{D}}_t \text{ are at positions } > 0, \\ \text{all holes in } \hat{\mathfrak{D}}_t \text{ are at positions } \leq 0 \end{array} \right). \quad (2)$$

- ▶ the l.h.s. of (2) is very close to  $\mathbb{P}(\mathfrak{H} < g(k_N, c))$
- ▶ the r.h.s. of (2) concerns two events involving only particles/holes
- ▶ using couplings, we can eventually compute the r.h.s. with the CLT for ASEP: It converges to  $F_{\text{GUE}}(cf(\alpha))$

Thank you for your attention !